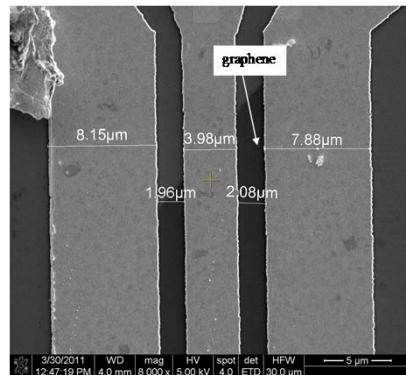


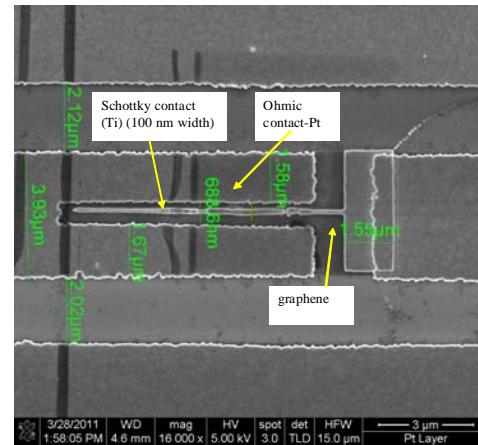
Nanoscale diodes without p-n junctions.

Mircea Dragoman
IMT –Bucharest

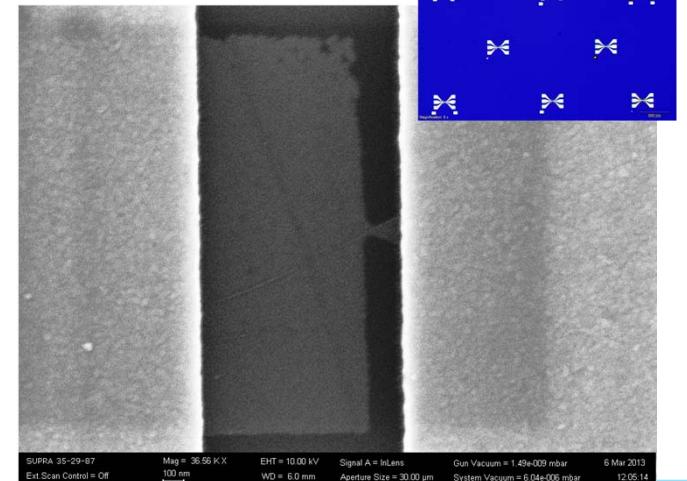
Coplanar electrodes on
graphene-graphene radio



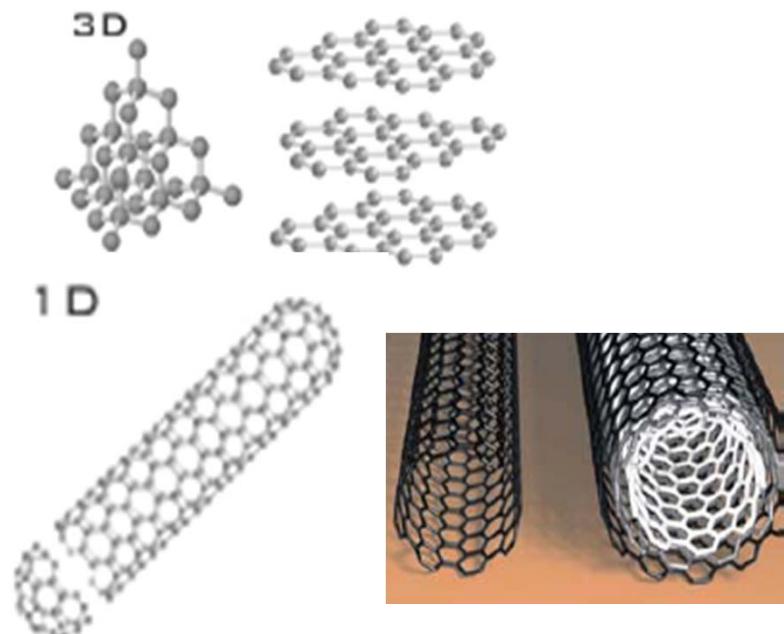
Dissimilar electrodes –
Schottky diode



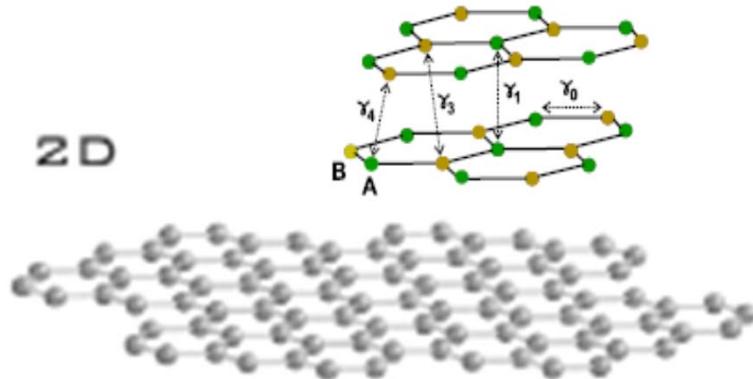
Geometric diodes-
graphene on wafer



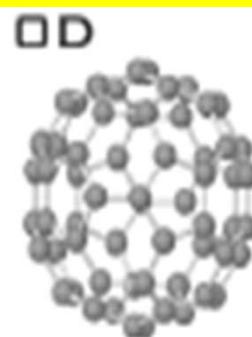
Carbon-based materials



S. Iijima, Nature 354 (1991) 56



K.S. Novoselov et al., Science 306 (2004)
666
Theory: P.R. Wallace, Phys. Rev. 71
(1947) 622



H.W. Kroto et al.,
Nature 318 (1985)
162



Graphene: Should it exist?

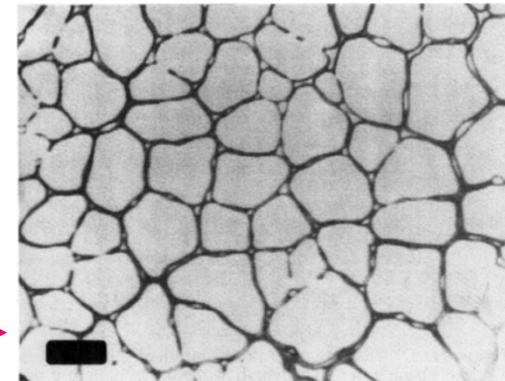
Two-dimensional crystals are thermodynamically unstable!

R.E. Peierls, Ann. I.H. Poincare 5 (1935) 177

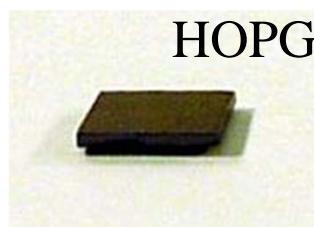
N.D. Mermin, Phys. Rev. 176 (1968) 250

L.D. Landau, E.M. Lifshitz, Statistical Physics, Pergamon (1980)

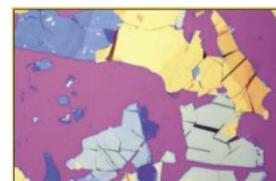
thickness < 20 nm



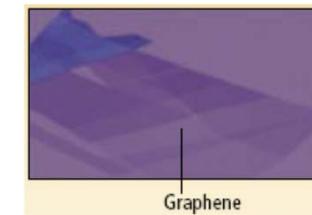
However.....



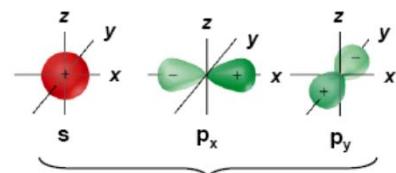
rub
deposit on
 $300 \mu\text{m SiO}_2$



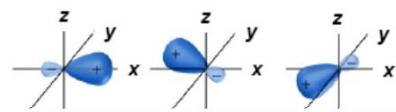
A.K. Geim, P. Kim, Sci. Am., April 2008, 90



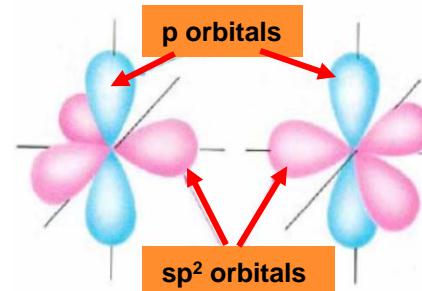
Why the current is flowing in graphene?



Isolated atom orbitals



Molecule-hybridization of orbitals; localized electrons



sp^2 carbon

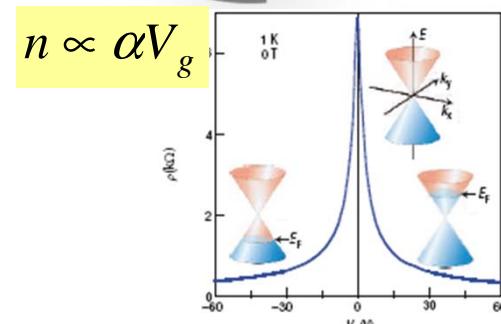
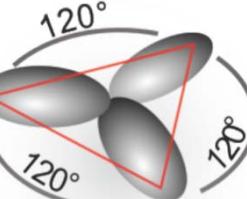
sp^2 carbon

Current flow



Carbon-carbon double bond

σ -electrons



K.S. Novoselov et al., Science 315 (2007) 1379

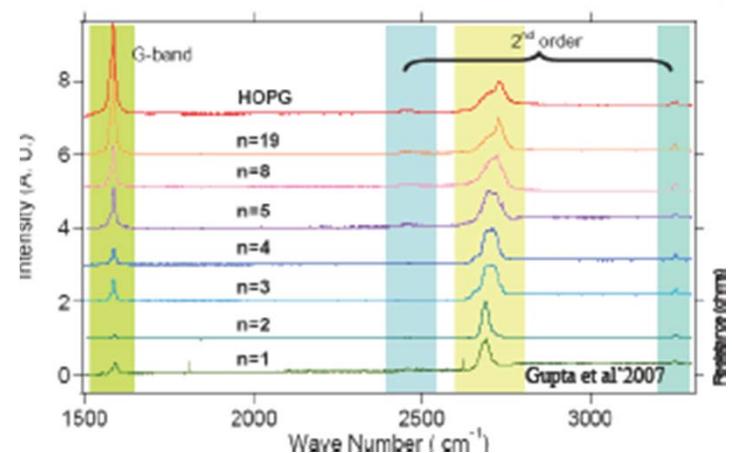


How we can identify graphene ?

Optical view of different graphene layers (mono, bi and multilayers)



Z.H. Ni et. al., Nano Lett.,
pp.2758-2763 (2007)



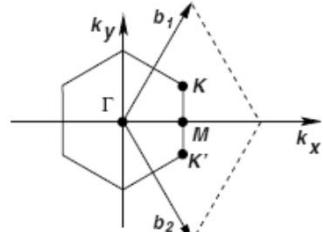
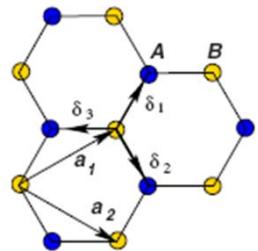
Raman
signature

Number of
layers

A.C. Ferrari et. al.,
Pyhs. Rev. Lett. 97 ,
187401 (2006)

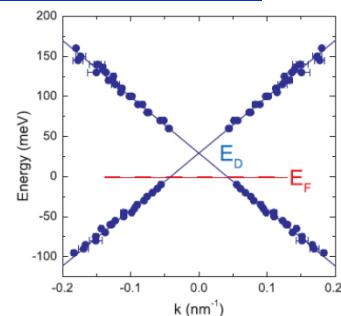


Band structure



$$E_{\pm}(k) = \pm \hbar v_F |k|$$

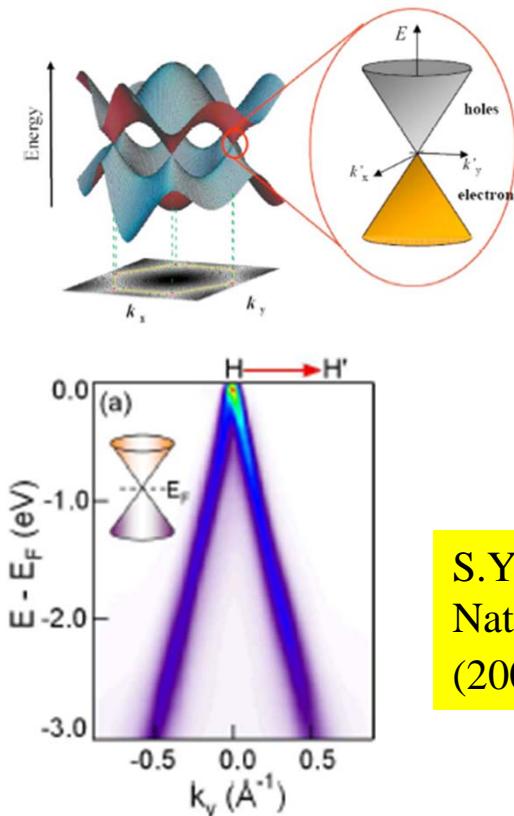
$$v_F \cong c / 300$$



D.L. Miller et al.,
Science 324 (2009)
924

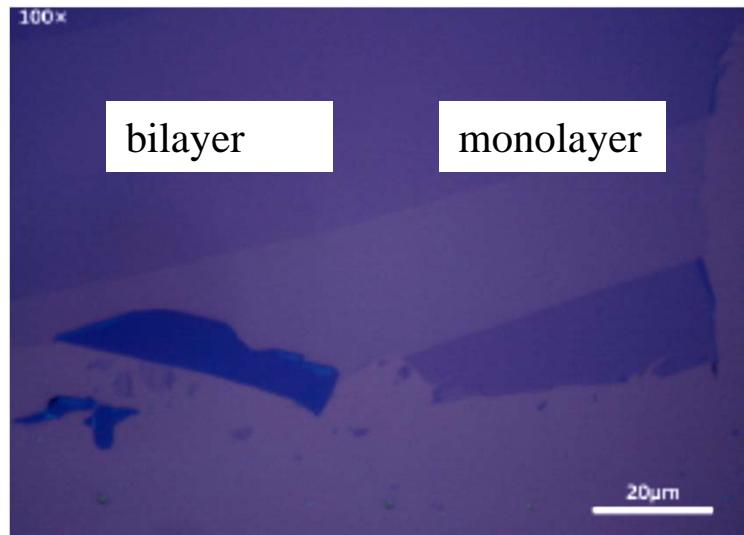
$$E = \sqrt{m^2 c^4 + p^2 c^2}$$

$$m = 0, \quad c = v_F$$



S.Y.Zhou et al.,
Nature Physics 2
(2006) 595

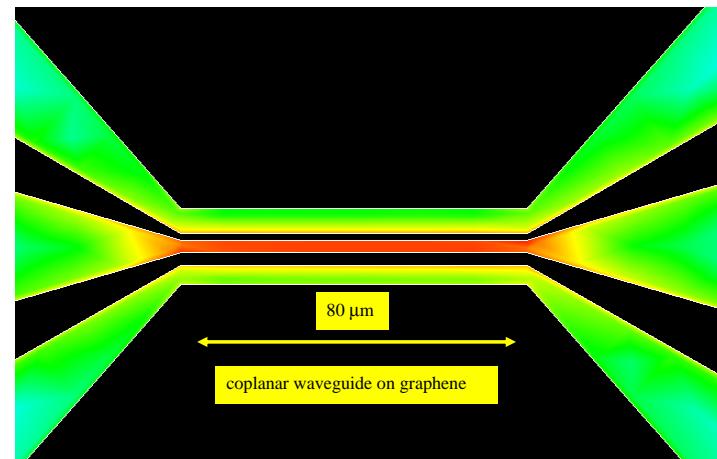
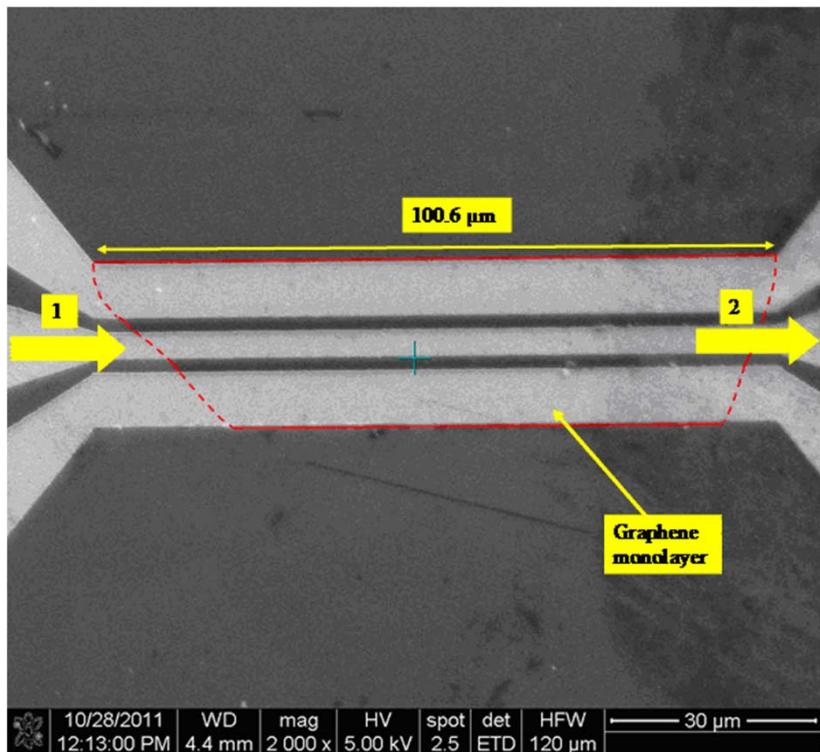
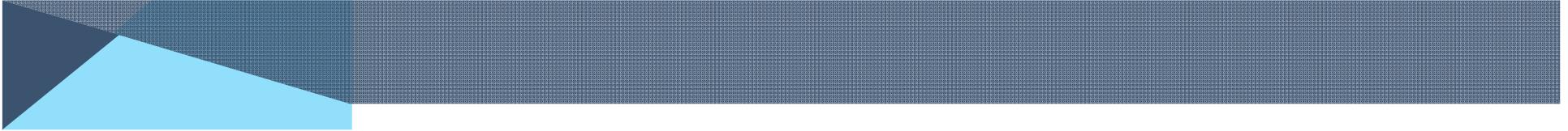




BALLISTIC
TRANSPORT

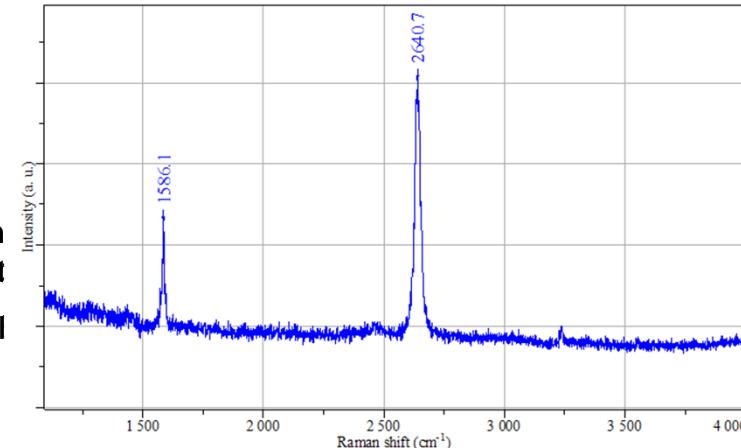
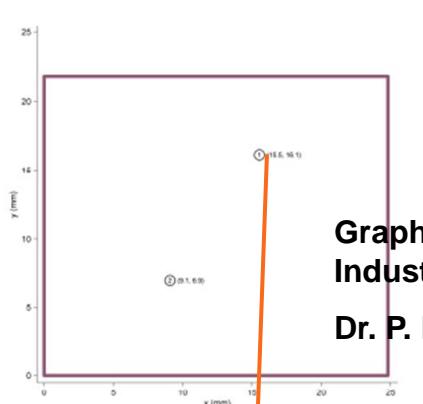
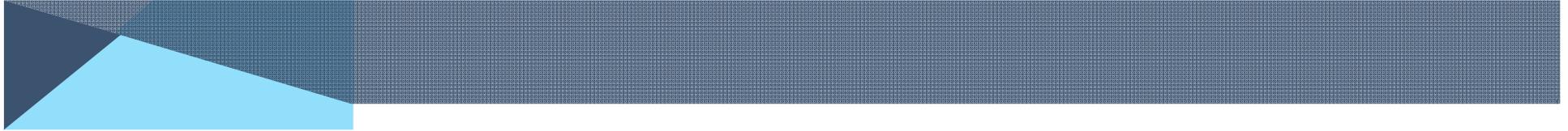
Parameter	Value and units	Observations
Thermal conductivity	5000 W/mK	Better thermal conductivity than in most crystals
Young modulus	1.5 TPa	Ten times greater than in steel
Mobility	40 000 cm ² V ⁻¹ s ⁻¹	At room temperature (intrinsic mobility) maximum mobility : 200 000 cm ² V ⁻¹ s ⁻¹) on suspended graphene or graphene on hexagonal BN substrate
Mean free path (ballistic transport)	=400 nm	At room temperature, but exceeds 1 μm in graphene on hexagonal BN substrate at room temperature
Electron effective mass	0	At room temperature
Hole effective mass	0	At room temperature
Fermi velocity	c/300=1000000 m/s	At room temperature



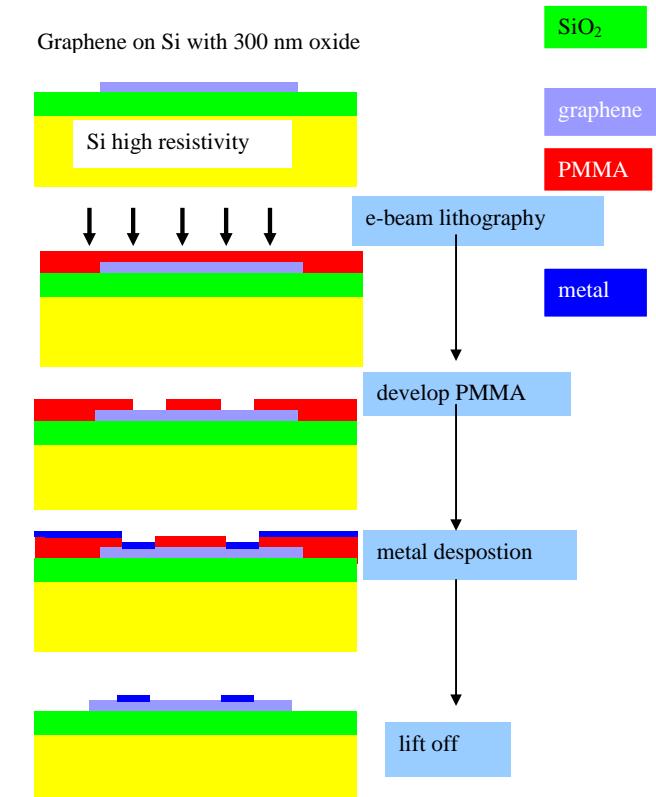


Simulation done at
100 GHz

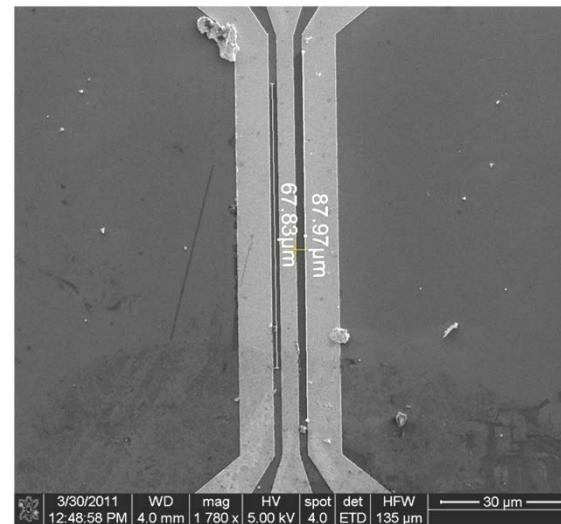
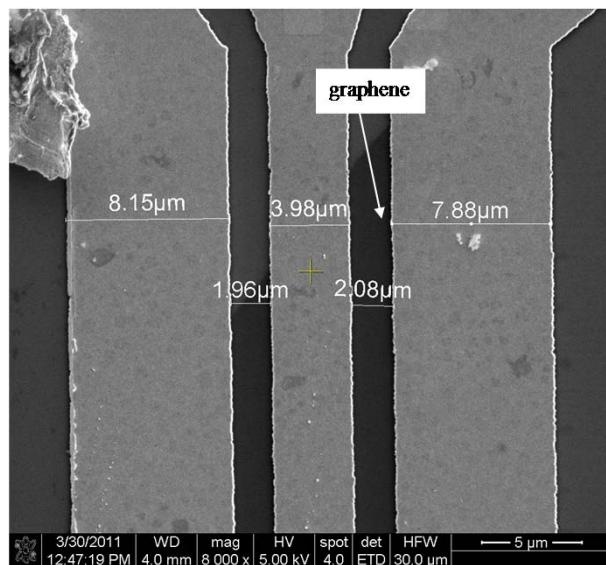


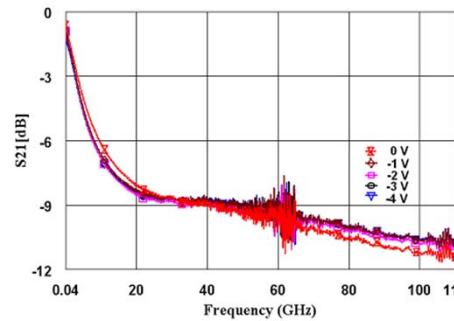
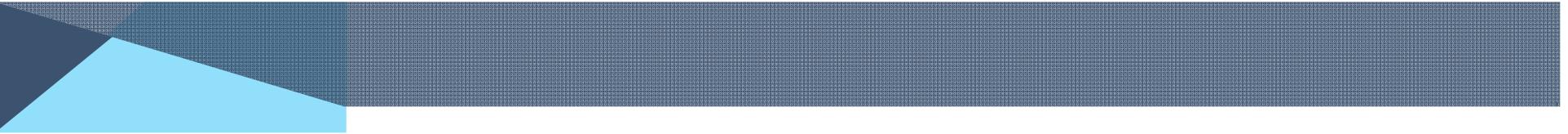


The ratio between the intensities of 2D and G peaks is around 2 telling us that we have a single layer graphene. The 2D and G Raman peaks positions are 2640 cm^{-1} and 1586cm^{-1} , respectively

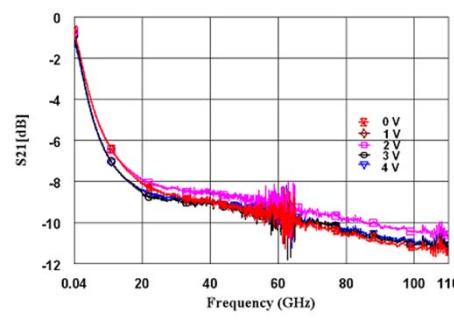


A CLOSER LOOK VIA SEM





Negative DC voltages

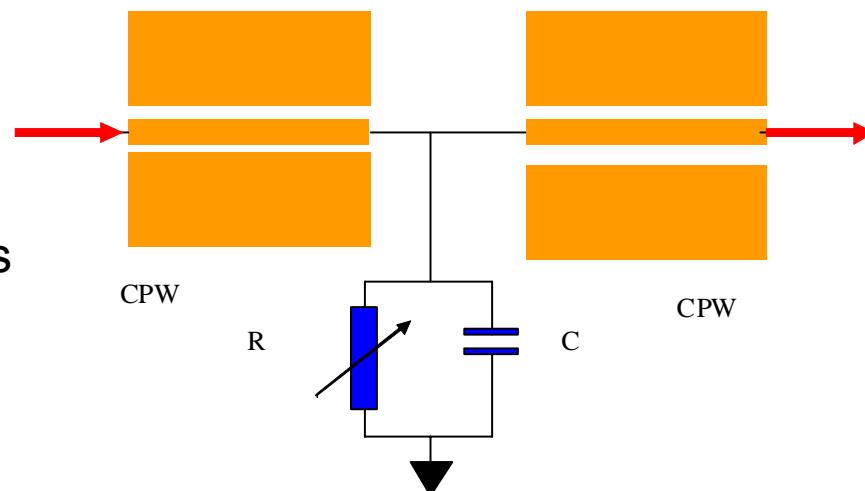


Positive DC voltages

(a)

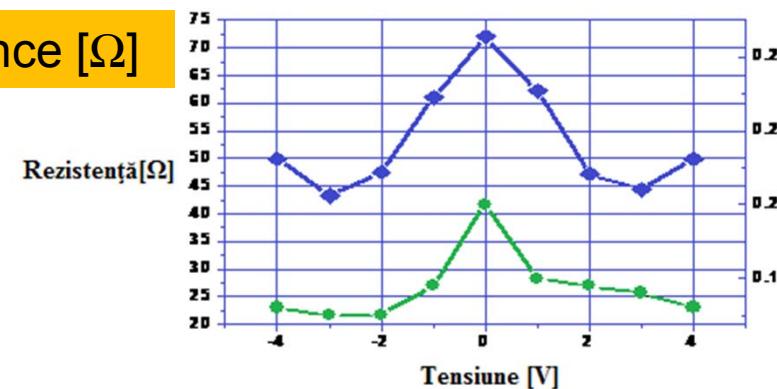
(b)

A SIMPLE MODEL



V_{bias} [V]	-4	-3	-2	-1	0	1	2	3	4
R[Ω]	49.7	42.9	46.8	61.75	73	63	46.6	43.6	49
C[pF]	0.173	0.17	0.174	0.179	0.205	0.181	0.179	0.177	0.173

Resistance [Ω]



Capacitance [pF]



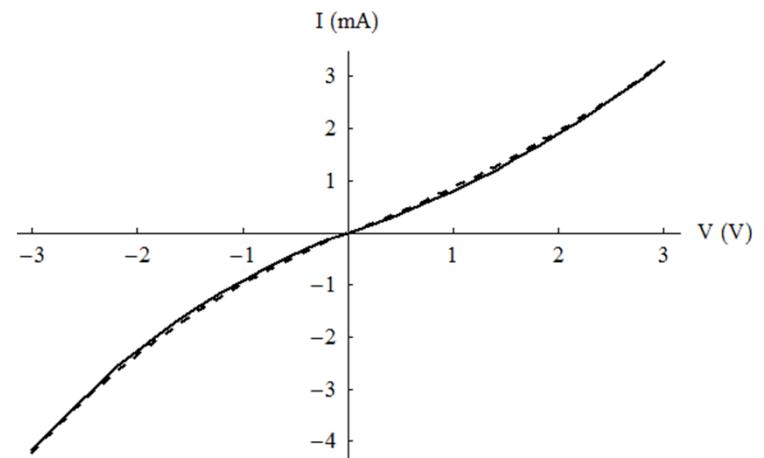
$$I = I_0[\exp(V / V_0) - 1] \quad (1)$$

I_0 and V_0 have the values 3.65 mA and 4.68 V for the positive polarization and -2.6 mA and -3.12 V for the negative polarization regime, respectively. Slightly asymmetric characteristics are typical in graphene devices and are due to graphene-substrate (in our case to graphene-CPW as well) interactions.

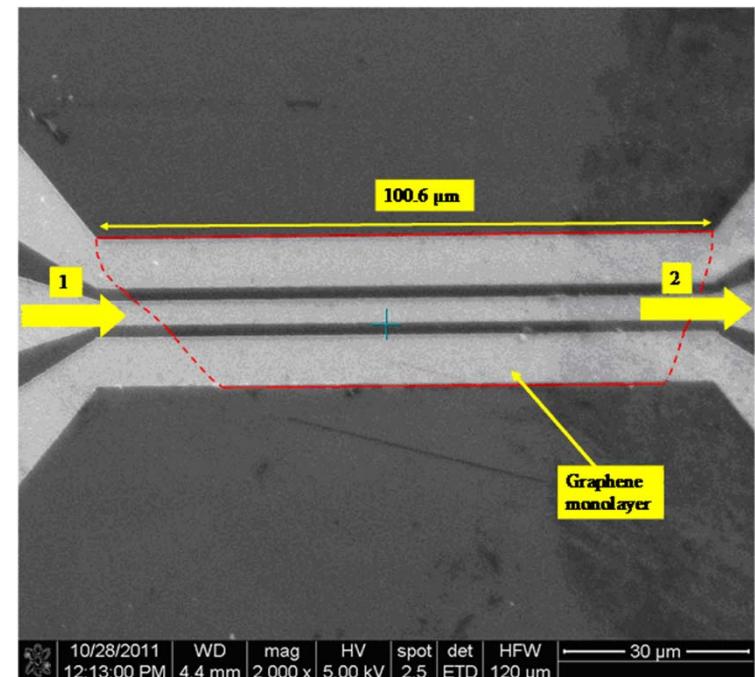
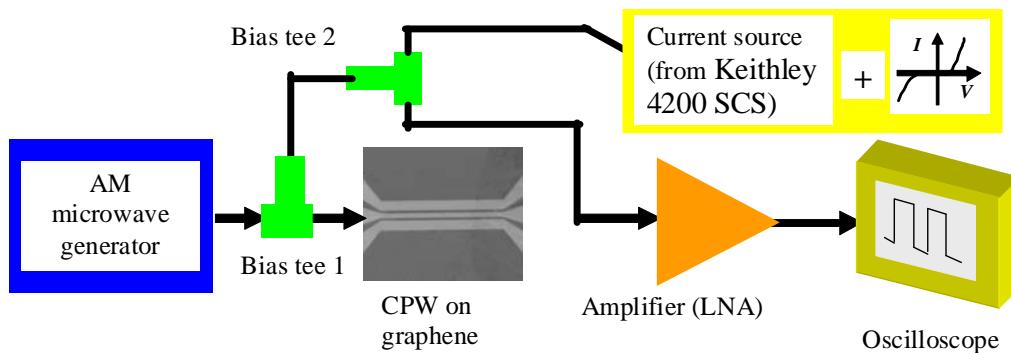
Choosing a operating point I_{av} and V_{av} and developing (1) in a Taylor series, around an operating point it results the demodulating term arround (I_{av} , V_{av}):

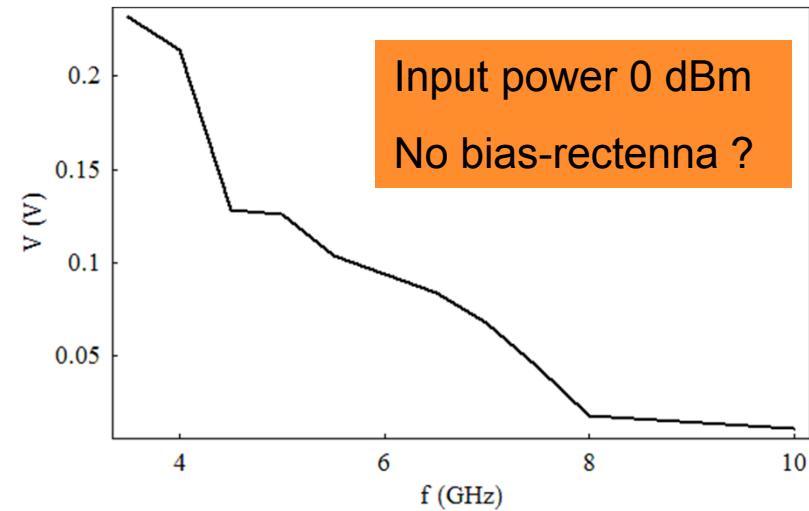
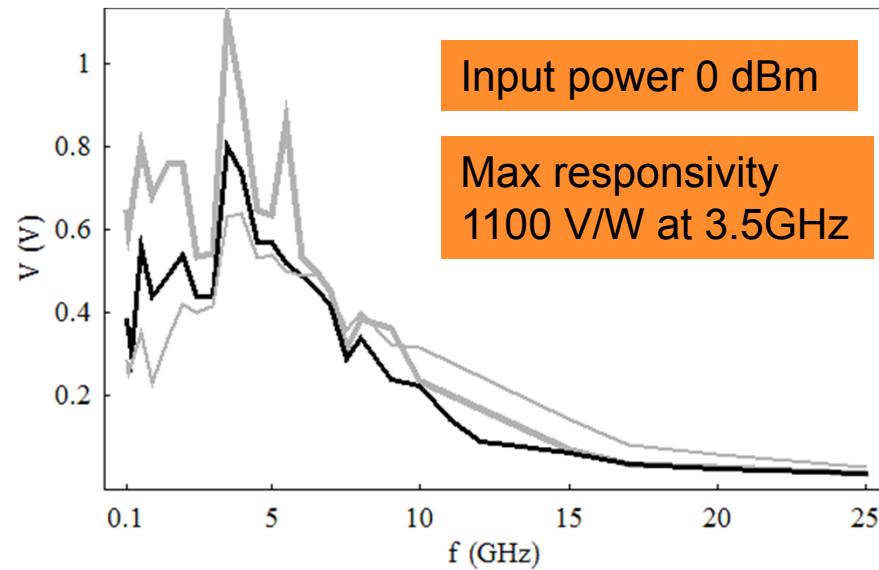
$$\Delta I = I - I_{av} = I_0 \frac{V_{RF}^2}{4V_0^2} \exp(V_{av} / V_0) \quad (2)$$

V_{RF} -the value of the RF signal

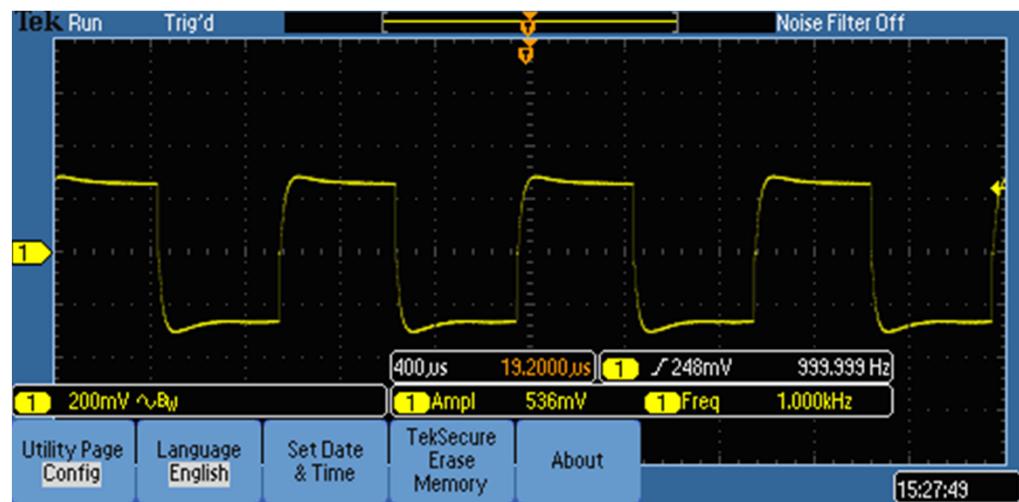


GRAPHENE RADIO





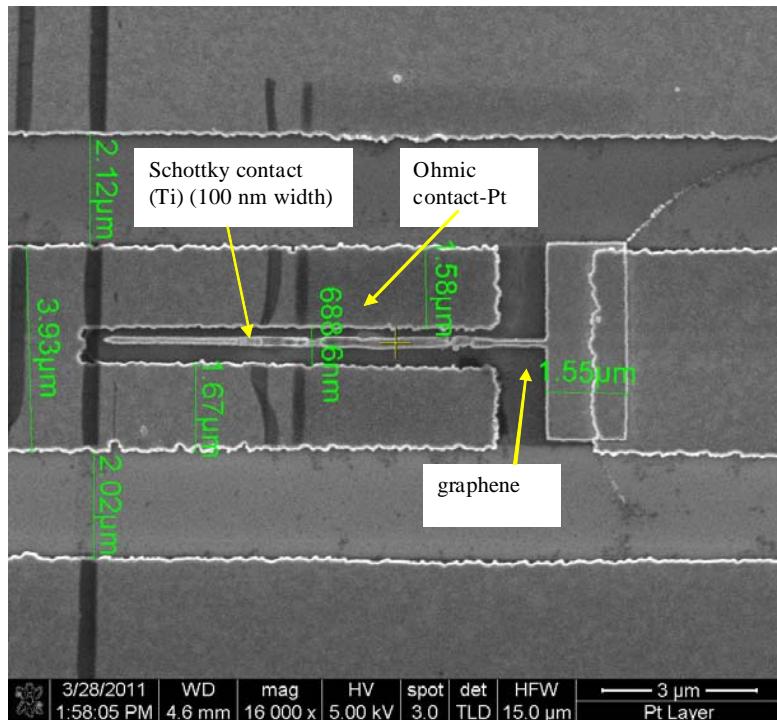
The detected DC voltage as a function of frequency for various DC currents: 1 mA (thin gray line), 2 mA (black line), 3 mA (thick gray line).



Demodulated signal
in time 1 kHz



SCHOTTKY DIODE VIA DISSIMILAR METALS



Schottky metals for graphene

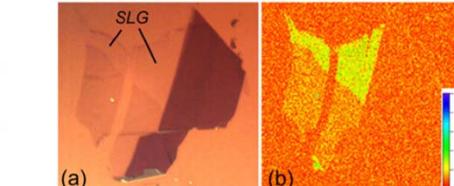
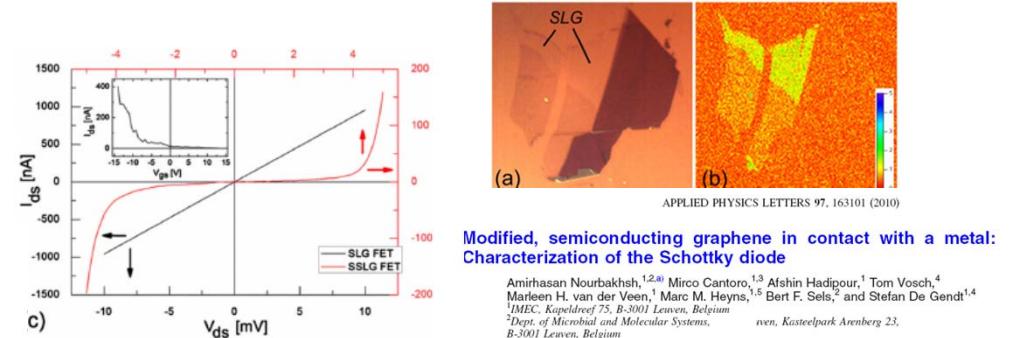
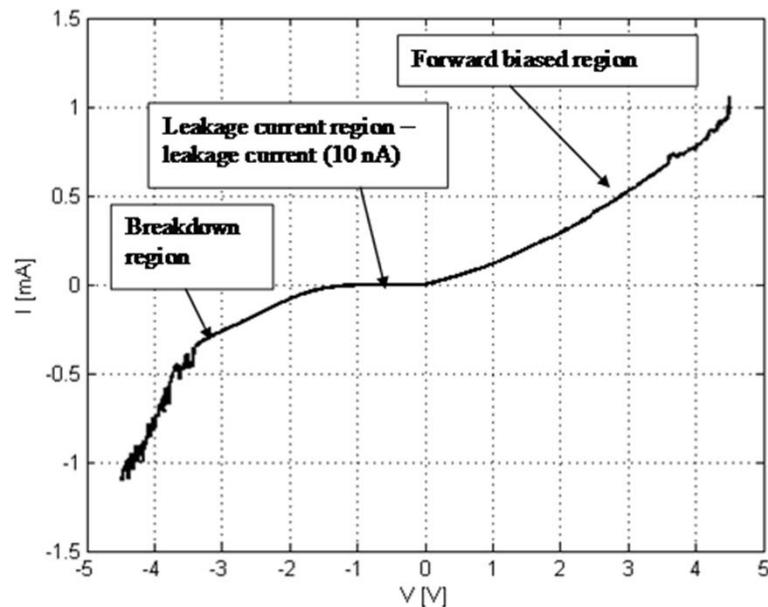
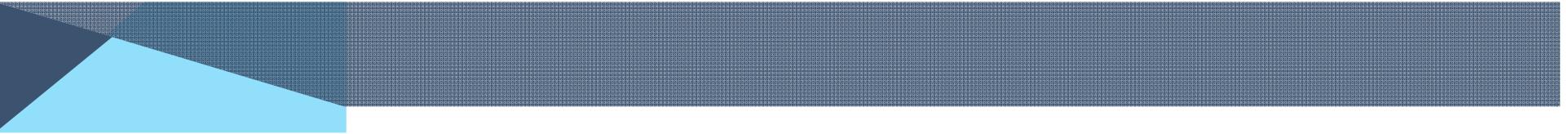
Metal	Work function (eV)
Al	-4.27 eV (the best)
Cr	-4.5 eV
Ti	-4.33 eV

Graphene work function -4.5 eV

Ohmic metals for graphene

Metal	Work function (eV)
Pd	-5.12 eV
Pt	-5.6 eV





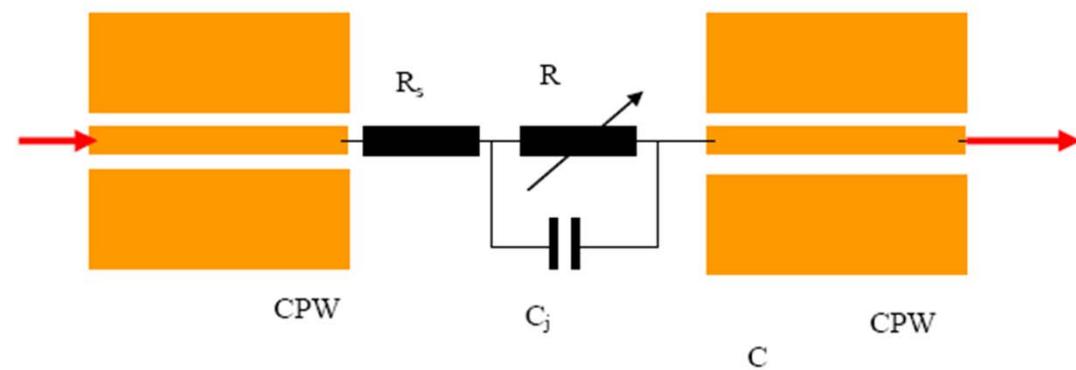
APPLIED PHYSICS LETTERS 97, 163101 (2010)

Modified, semiconducting graphene in contact with a metal: Characterization of the Schottky diode

Amirhasan Nourbakhsh,^{1,2,a)} Mirco Cantoro,^{1,3} Afshin Hadipour,¹ Tom Vosch,⁴ Marleen H. van der Veen,¹ Marc M. Heyns,^{1,5} Bert F. Sels,² and Stefan De Gendt^{1,4}
¹IMEC, Kapeldreef 75, B-3001 Leuven, Belgium
²Dept. of Microbial and Molecular Systems, B-3001 Leuven, Belgium
³Dept. of Physics and Astronomy, K. U. Leuven, Celestijnenlaan 200d, B-3001 Leuven, Belgium
⁴Dept. of Chemistry, K. U. Leuven, Celestijnenlaan 200d, B-3001 Leuven, Belgium
⁵Dept. of Metallurgy and Materials Engineering, B-3001 Leuven, Belgium

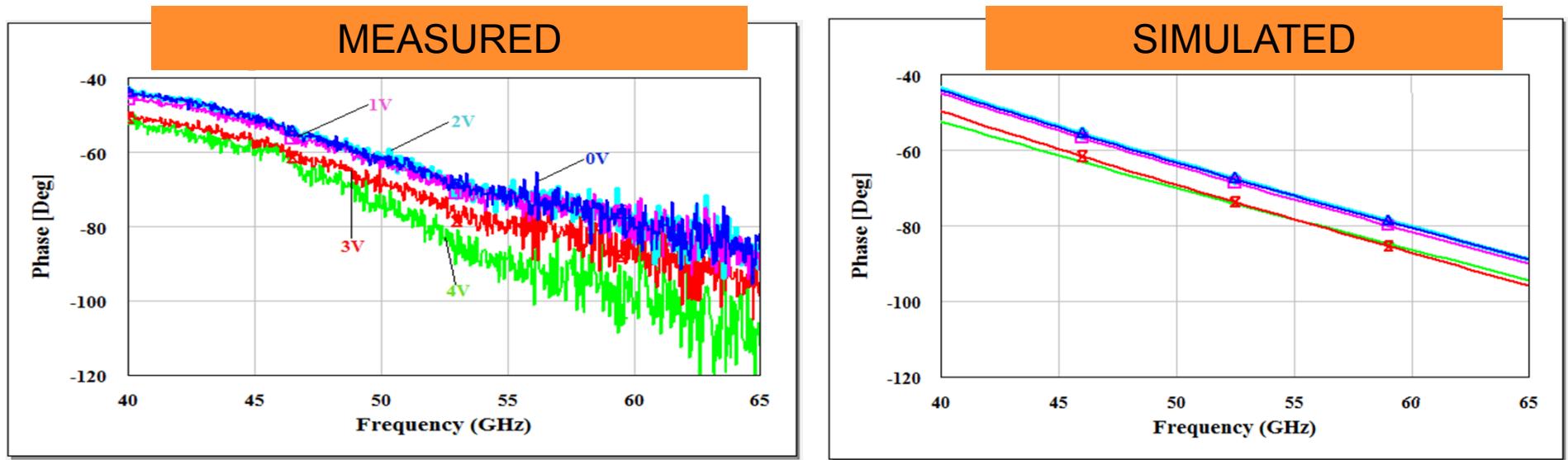
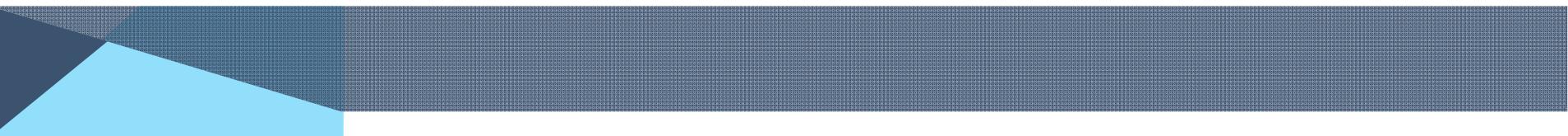
Our results-currents
at mA level!





Bias voltage (V)	R_s [Ω]	R_J [$k\Omega$]	C_J [fF]
0V	60	12	3.5
1V	60	8	3.5
2V	60	8	3.5
3V	60	0.85	3.5
4V	60	0.61	3.5

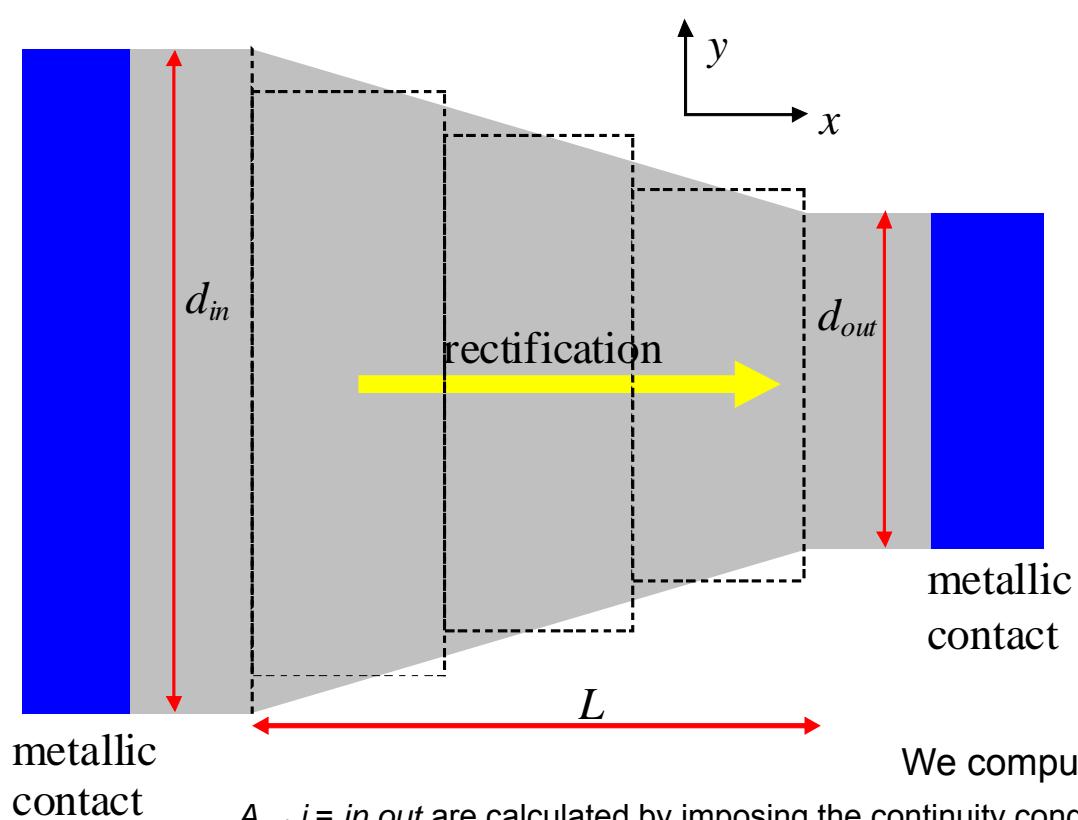




A GRAPHENE PHASE SHIFTER



BALLISTIC GEOMETRIC GRAPHENE DIODE



$$d_{in} > d_{out}$$

$$D_{out}=100\text{nm}$$

$$D_{in}=20 \text{ nm}$$

$$L=200\text{nm}$$

$$d_j = d_{in} - (d_{in} - d_{out}) j / (N + 1)$$

$$V_j = -jeV / (N + 1) \quad \text{Potential energy}$$

$$k_{n,j} = \text{sgn}(E - V_j) \sqrt{(E - V_j)^2 / (\hbar^2 v_F^2) - (n\pi/d_j)^2}$$

$$T = \sum_{n=1}^{N_{out}} |A_{n,out}|^2 / \sum_{n=1}^{N_{in}} |A_{n,in}|^2$$

We compute:

$A_{n,j}$ $j = in, out$ are calculated by imposing the continuity conditions at each interface for the spinorial solutions of the Dirac equation in region j :

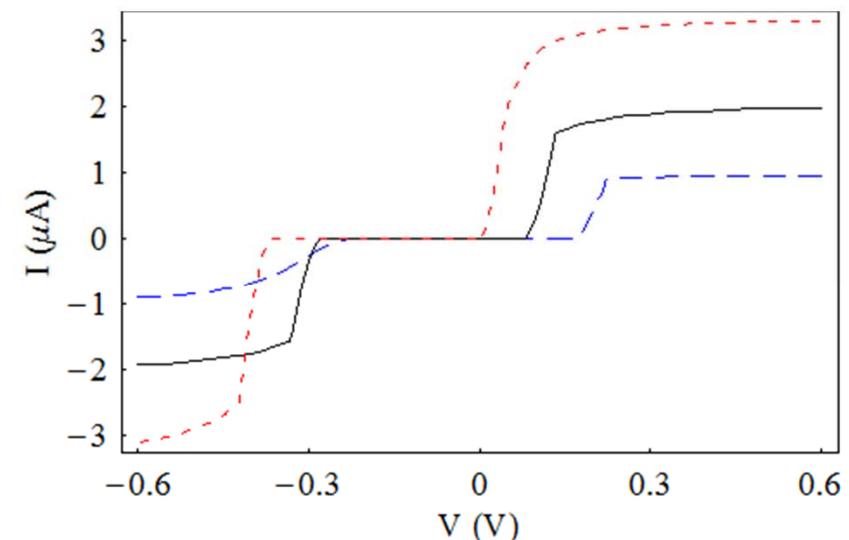
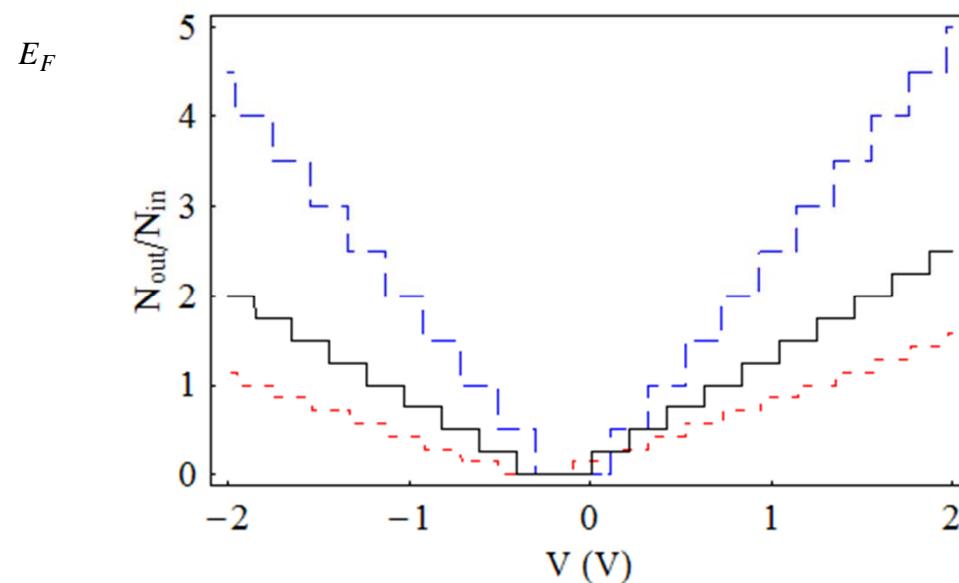
$$\Psi_j(x, y) = \begin{pmatrix} \sum_{n=1}^{N_j} [A_{n,j} \exp(ik_{n,j}x) + B_{n,j} \exp(-ik_{n,j}x)] \sin(2n\pi y/d_j) \\ \sum_{n=1}^{N_j} [A_{n,j} \exp(ik_{n,j}x) - B_{n,j} \exp(-ik_{n,j}x)] \sin(2n\pi y/d_j) \end{pmatrix}$$



$$N_j = \text{Int}[d_j |E - V_j| / (\pi \hbar v_F)]$$

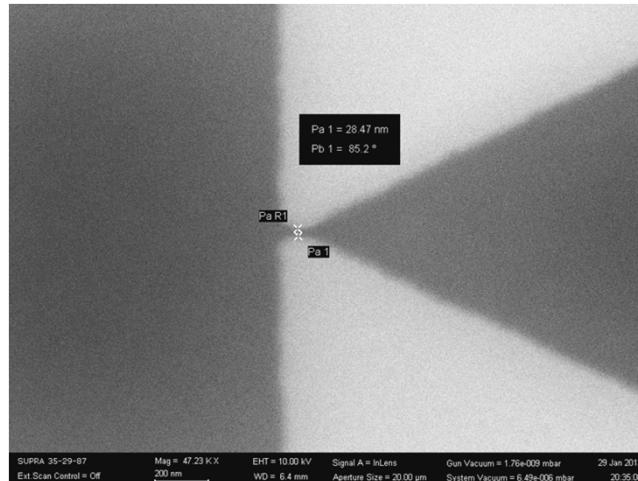
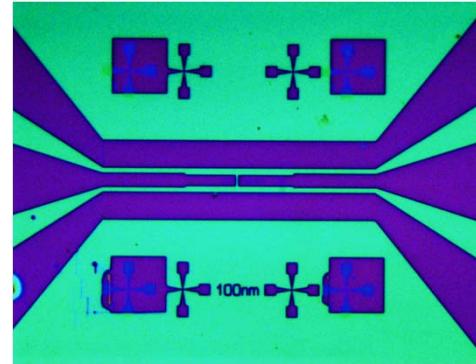
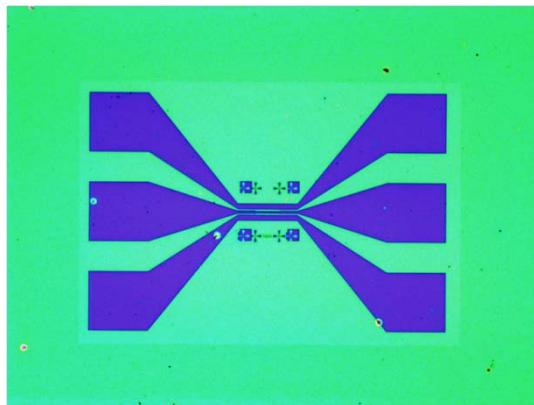
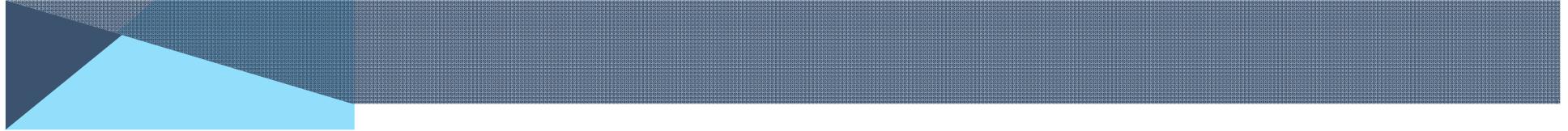
Number of modes

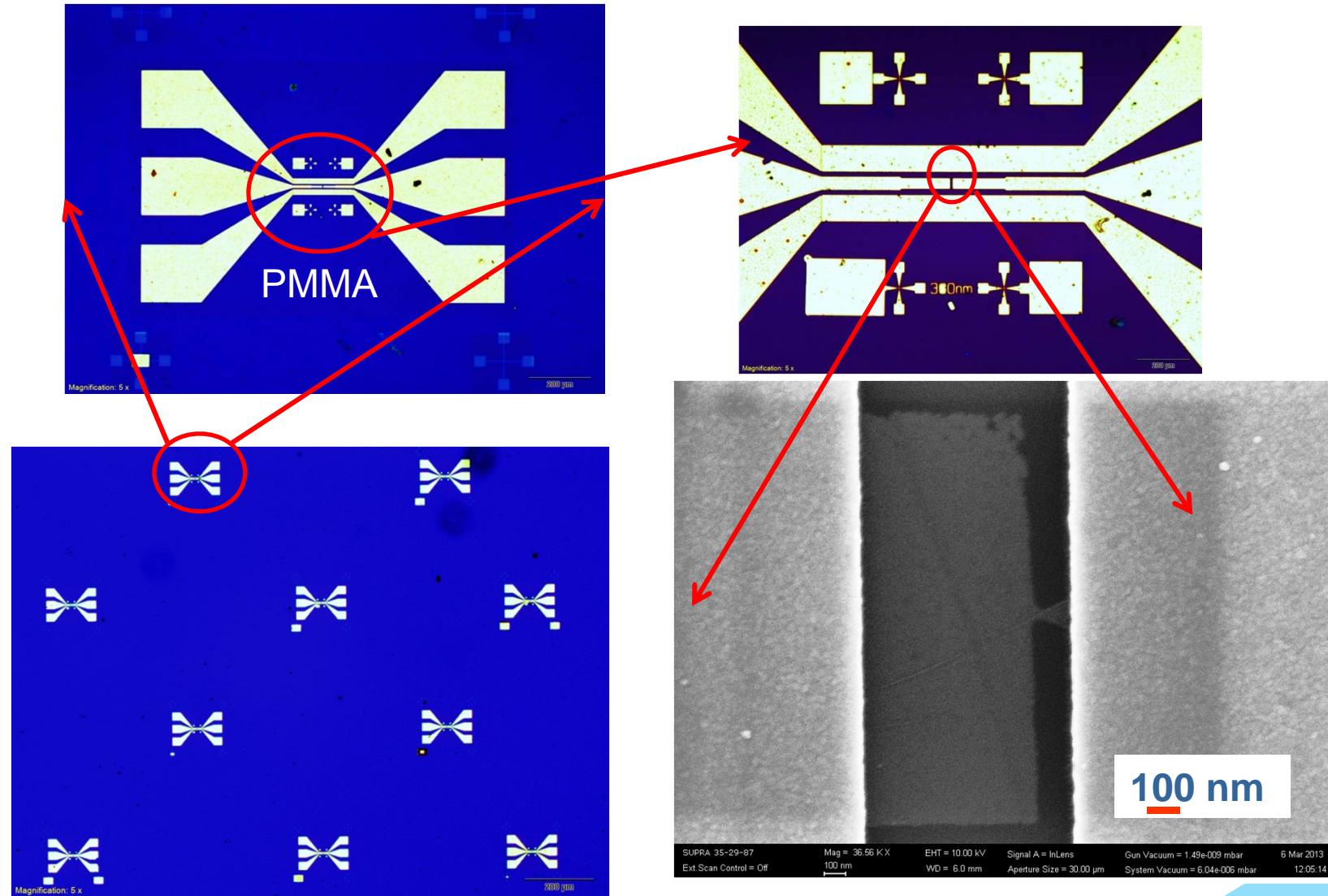
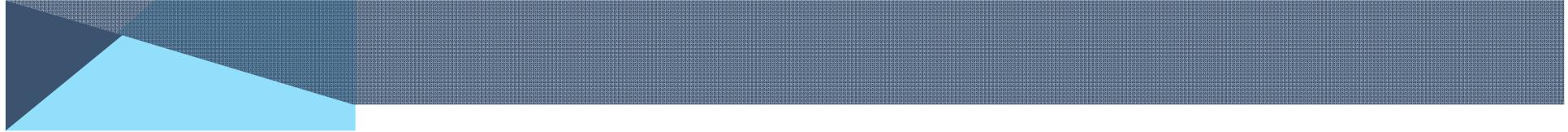
The results are independent on the number of discretization regions N . Although there are a finite number of outgoing modes for both voltage polarizations, in all cases there is a voltage range in which no charge carriers are transmitted since for these V values the number of outgoing modes, and hence the current, vanishes. This region, with a width given by $\pi \hbar v_F / d_{out}$



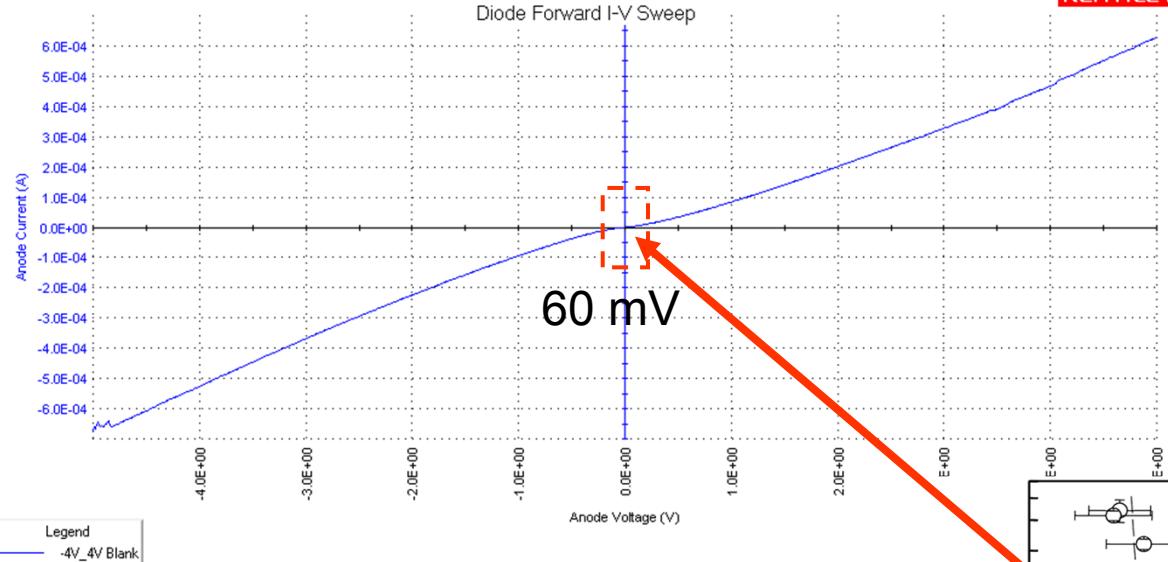
$E_F = 0$ (blue dashed line), 0.1 eV (solid black line) and 0.2 eV (red dotted line).







03/08/2013 16:01:05

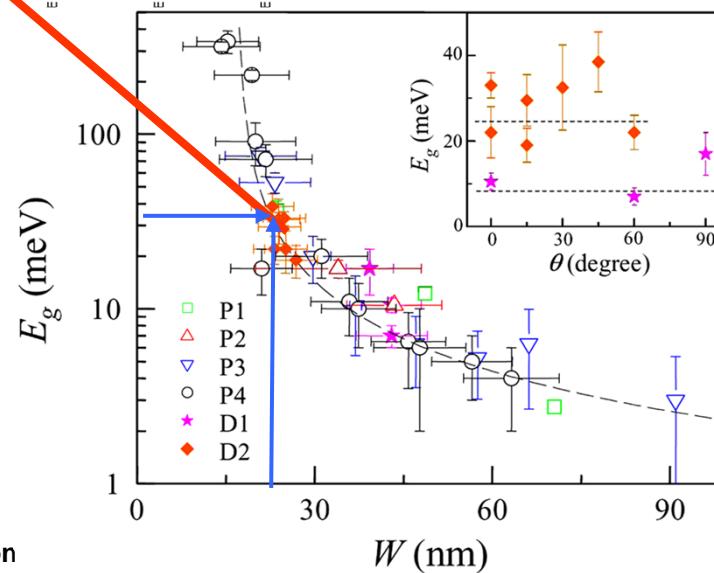


$k_b T = 27 \text{ meV}$ at
room
temperature

$E_g = 60 \text{ meV}$

$C = 3 \text{ aF}$
 $R = 10K \Omega$
 $\tau = 30\text{fs}$ $f_c = 6 \text{ THz}$

Phys. Rev. Lett. 98, 206805 (2007)
Energy Band-Gap Engineering of Graphene Nanoribbon



ACKNOWLEDGEMENTS

I would like to thank to many scientists which helped me in the quest in the area of carbon nanoelectronics, most of them being co-authors of many papers published in the last period of time. **My wife Daniela** has helped me with her deep knowledge in the area of quantum mechanics ,solid state physics and graphene physics . **My colleagues from IMT Bucharest (Alex Muller, Dan Neculoiu, Adrian Dinescu, Alina Cismaru, Antonio Radoi,), dr. George Konstantinidis from FORTH Heraklion, dr. George Deligeorgis from LAAS Toulouse** were behind almost any device reported in this talk especially due to their invaluable knowledge in the area of semiconductor technology. I am grateful **Prof. Hans Hartnagel** who during many years has helped and guided me in the area of microwaves and nanoelectronics.

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